

Problem Set 4

Theoretical Solid State Physics (SoSe2017)

Due: Due Thursday, May 18, 2017; at the beginning of class

Problem 1: Berry phase – an example

Consider a spin 1/2 in an adiabatically varying Zeeman field $\mathbf{B}(t)$, as described by the Hamiltonian

$$H = -\mathbf{B}(t) \cdot \sigma, \quad (1)$$

where σ denotes the vector of Pauli matrices. Parametrize the magnetic field as

$$\mathbf{B} = B(\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta). \quad (2)$$

- Find the eigenvectors of the Hamiltonian for an arbitrary, but fixed magnetic field. Make sure that your result is single valued as function of the magnetic field.
- Use your result in (a) to compute the corresponding Berry connection. Relate your result explicitly to the vector potential for a magnetic monopole which we encountered earlier.
- Compute the corresponding Berry curvature and show that the Berry phase accumulated when the magnetic field varies in a closed loop is equal to half the solid angle subtended by the magnetic field loop.

Problem 2: Shortcuts to adiabaticity

Shortcuts to adiabaticity are currently a popular topic in research and address the following question: Imagine that you find the adiabatic dynamics of some Hamiltonian $H_0(t)$ attractive and want to implement it (e.g., because there are no transitions between instantaneous eigenstates $|\psi_n(t)\rangle$ with instantaneous eigenenergies $E_n(t)$), but you do not have the patience to wait long enough so that the dynamics is really adiabatic. Then you would like to find a way to implement the adiabatic dynamics in a finite (and possibly short) time. One possibility is to search for a modified Hamiltonian $H(t) = H_0(t) + H_1(t)$ which implements the desired adiabatic dynamics as its *exact* quantum dynamics. It turns out that this problem can be solved quite generally.

- Show that the time evolution operator describing the adiabatic dynamics of $H_0(t)$ can be written as

$$\mathcal{U}(t) = \sum_n e^{-i \int_0^t dt' E_n(t') + i \gamma_n(t)} |\psi_n(t)\rangle \langle \psi_n(0)|, \quad (3)$$

where $\gamma_n(t) = i \int_0^t dt' \langle \psi_n(t') | \partial_{t'} \psi_n(t') \rangle$ denotes the Berry phase.

- Show that given a time evolution operator $\mathcal{U}(t)$, one can find the corresponding Hamiltonian for which $\mathcal{U}(t)$ describes the exact quantum dynamics through

$$H(t) = i[\partial_t \mathcal{U}] \mathcal{U}^\dagger. \quad (4)$$

- Now insert the desired adiabatic time evolution operator from (a) into the general expression in (b) to obtain

$$H_1(t) = i \sum_n (|\partial_t \psi_n\rangle \langle \psi_n| - |\psi_n\rangle \langle \psi_n| \partial_t \psi_n \langle \psi_n|). \quad (5)$$

for the so-called counter-diabatic term $H_1(t)$.

- Explicitly evaluate the counter-diabatic term for the system which you treated in problem 1.

This version of shortcuts to adiabaticity is due to Michael Berry, J. Phys. A: Math. Theor., **42** 365303 (2009) (which you can turn to if you need help in solving this problem). Such protocols can actually be implemented experimentally, see, e.g., M.G. Bason *et al.*, Nat. Phys. **8**, 147 (2012).

Problem 3: Dimerized site energies

Solve for the band structure of the 1D tight-binding Hamiltonian

$$H = \sum_j \{(-1)^j \epsilon |j\rangle\langle j| - t[|j\rangle\langle j+1| + |j+1\rangle\langle j|]\} \quad (6)$$

with dimerized site energies $\pm\epsilon$. At half filling, is the model a metal or an insulator?