


Introduction to Finite Elements


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Seminar in Engineering Analysis

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Outline


- Review midterm
- Introduction to finite elements
- Basic approaches to finite elements
 - Will start material originally scheduled for April 22
 - Parallel reading for this week is pages 711 to 739 in Hoffman
- Example application in one dimension



Midterm Results

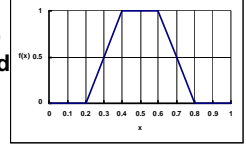
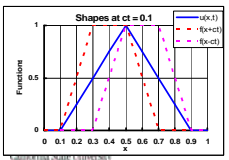
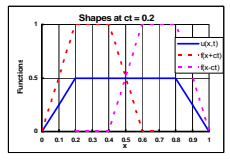

- Number of students = 12
- Maximum possible = 100
- Mean = 45.1
- Median = 43.5
- Standard deviation = 28.2
- Grade distribution

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|----|----|----|----|----|----|
| 2 | 10 | 14 | 26 | 35 | 40 |
| 47 | 65 | 71 | 73 | 75 | 83 |



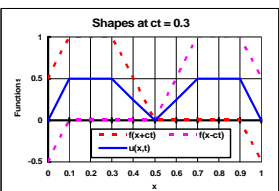

Problem One

- Sketch the wave equation solution, $u(x,t)$ at times $ct = 0.1, 0.2,$ and 0.3 for $f(x) = u(x,0)$ as shown at the right.
 - $u(x,t) = [f(x+ct) + f(x-ct)]/2$

Problem One II

- For $ct = 0.3$ must account for periodic extensions of initial conditions for $x < 0$ and $x > 1$ entering region $0 \leq x \leq 1$
 - These are inverse of original initial condition





Problem Two

- Outline solution for potential $u(x,y,t)$

$$\frac{\partial u}{\partial t} = \alpha \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) \quad 0 \leq x \leq L, 0 \leq y \leq H, t \geq 0 \quad u(x,y,0) = U_0$$

$$u(0,y,t) = 0 \quad u(L,y,t) = U_1 \quad u(x,0,t) = 0 \quad u(x,H,t) = U_2$$
- $u(x,y,t) = v(x,y,t) + w(x,y)$
 - $v(x,y,t)$ is diffusion equation solution with homogenous boundary conditions
 - $w(x,y)$ is Laplace equation solution requiring superposition: $w(x,y) = w_1(x,y) + w_2(x,y)$
 - $w_1(x,y)$ and $w_2(x,y)$ each have only one nonhomogenous boundary condition



Problem Two II

- Solutions in notes

$$v(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} C_{nm} e^{-\left[\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2\right] \alpha t} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi y}{H}\right)$$

$$w_1(x, y) = \frac{4U_2}{\pi} \sum_{n=0}^{\infty} \frac{\sin\left(\frac{(2n+1)\pi x}{L}\right) \sinh\left(\frac{(2n+1)\pi y}{L}\right)}{(2n+1) \sinh\left(\frac{(2n+1)\pi H}{L}\right)}$$

- Solution for $w_2(x,y)$ found by substituting x and y, L and H (and U_1 for U_2) in w_1 solution

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Problem Two II

$$w_2(x, y) = \frac{4U_1}{\pi} \sum_{n=0}^{\infty} \frac{\sin\left(\frac{(2n+1)\pi y}{H}\right) \sinh\left(\frac{(2n+1)\pi x}{H}\right)}{(2n+1) \sinh\left(\frac{(2n+1)\pi L}{H}\right)}$$

- Need to use solutions for $w(x,y)$ to find coefficients in $v(x,y,t)$ solution

$$C_{pq} = \frac{4}{HL} \int_0^H \int_0^L [U_0 - w_1(x, y) - w_2(x, y)] \sin\left(\frac{p\pi x}{L}\right) \sin\left(\frac{q\pi y}{H}\right) dx dy$$

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Problem Three

- Replace each f_k in equation below by its Taylor series to order of the error

$$f'_i = \frac{-f_{i+2} + 8f_{i+1} - 8f_{i-1} + f_{i-2}}{12h}$$

$$f_{i+2} = f_i + f'_i 2h + \frac{f''_i (2h)^2}{2!} + \frac{f'''_i (2h)^3}{3!} + \frac{f^{(4)}_i (2h)^4}{4!} + \frac{f^{(5)}_i (2h)^5}{5!} + \dots$$

$$f_{i+1} = f_i + f'_i h + \frac{f''_i (h)^2}{2!} + \frac{f'''_i (h)^3}{3!} + \frac{f^{(4)}_i (h)^4}{4!} + \frac{f^{(5)}_i (h)^5}{5!} + \dots$$

$$f_{i-1} = f_i - f'_i h + \frac{f''_i (-h)^2}{2!} + \frac{f'''_i (-h)^3}{3!} + \frac{f^{(4)}_i (-h)^4}{4!} + \frac{f^{(5)}_i (-h)^5}{5!} + \dots$$

$$f_{i-2} = f_i + f'_i (-2h) + \frac{f''_i (-2h)^2}{2!} + \frac{f'''_i (-2h)^3}{3!} + \frac{f^{(4)}_i (-2h)^4}{4!} + \frac{f^{(5)}_i (-2h)^5}{5!} + \dots$$

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Problem Three II

- After the algebra

$$\frac{-f_{i+2} + 8f_{i+1} - 8f_{i-1} + f_{i-2}}{12h} = \frac{(-2+8+8-2)f'_i h}{12h} + \frac{(-32+8+8-32)f''_i h^2}{5!(12h)} + \dots = f'_i + O(h^4)$$

- Get first derivative of e^x at $x = 0$ for $h = 0.1$ and $h = 0.05$, and find the order of the error

$$f'_i = \frac{-e^{0+2(0.1)} + 8e^{0+0.1} - 8e^{0-0.1} + e^{0-2(0.1)}}{12(0.1)} = 0.999996662696$$

$$f'_i = \frac{-e^{0+2(0.05)} + 8e^{0+0.05} - 8e^{0-0.05} + e^{0-2(0.05)}}{12(0.05)} = 0.99999979160465$$

$$n \approx \frac{\log(\epsilon_2) - \log(\epsilon_1)}{\log(h_2) - \log(h_1)} = \frac{\log(2.08395 \times 10^{-7}) - \log(3.3373 \times 10^{-6})}{\log(0.05) - \log(0.1)} = 4.001$$

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Why Finite Elements

- A different approach for converting differential equations into algebraic equations
 - More complicated approach than finite differences
 - Capable of handling complex geometry more accurately than finite differences
 - Will not cover finite-volume a melding of the two
 - Uses concept of interpolation polynomials
 - Introduces gradients naturally

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Finite Element Methods

- Designed for 2D and 3D geometries
- Can use for 1D case as example
- Basic idea is to divide region into small elements (line, area, volume)
- Use interpolating polynomial for each element
 - Represent both geometry (independent variables) and dependent variable
 - Polynomials called basis functions or shape functions

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Finite Element Methods II

- Start analysis for region
- Look at set of small elements in region
- Assemble analyses for individual elements into a set of nodal equations for the entire region
 - Result is set of algebraic equations for the dependent variable at nodes that are points on elements
 - Converts differential equations into a set of algebraic equations at distinct points

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Finite Element Analyses

- How do we represent differential equation in terms of polynomials?
 - Variational approach (Rayleigh-Ritz) method formulates problem as the maximum (or minimum) of a function integral
 - Apply directly to problems in solid mechanics governed by a variational principle (Hamilton's principle)
 - Method of weighted residuals
 - Applies to general differential equations
 - Least squares approaches

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Method of Weighted Residuals

- MWR uses shape functions $\phi_i(\mathbf{x})$ or $N_i(\mathbf{x})$ for each node in region to write approximate solution for true $u(\mathbf{x})$

$$\hat{u} = \sum_{i=1}^N u_i \phi_i(\mathbf{x})$$
- Have to find all the u_i values
- Write differential equation as $L(u) - b = 0$
- Weighting functions $w_i(\mathbf{x})$ for individual nodes in region that satisfy this equation

$$\int_{\Omega} w_i [L(\hat{u}) - b] d\Omega = 0 \quad i = 1, \dots, N$$

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MWR Example

- Poisson's equation for $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = -Q$
 $0 \leq x \leq L$ and $0 \leq y \leq H$
- Pick N points in region including nodes on the boundary
- Pick a set of weighting functions w_i for $L(u)$
 - Choice for w_i determines final form of algebraic equations solved
$$\int_0^L \int_0^H w_i \left[\frac{\partial^2 \hat{u}}{\partial x^2} + \frac{\partial^2 \hat{u}}{\partial y^2} + Q \right] dy dx = 0 \quad i = 0, \dots, N$$

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What are Weight Functions?

- Different approaches to MWR use different weight functions
- Galerkin's method uses $w_i = \phi_i$
 - Gives same result as variational approach when a variational principle is possible
- Collocation method uses $w_i = \delta(\mathbf{x} - \mathbf{x}_i)$
- Dirac delta function $\delta(\mathbf{x} - \mathbf{x}_i)$ result

$$\int_{\Omega} f(\mathbf{x}) \delta(\mathbf{x} - \mathbf{x}_i) d\Omega = f(\mathbf{x}_i)$$

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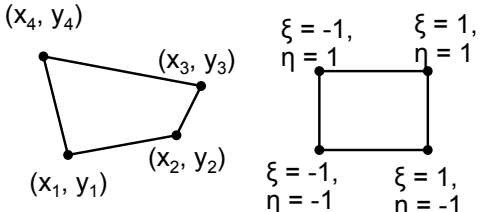
Shape (Basis) Functions

- Simplest shape functions are linear for 1D or bilinear for 2D
- For a linear element between nodes i (at $\xi = -1$) and $i + 1$ (at $\xi = 1$) we have $\phi_i = (1 - \xi)/2$ and $\phi_{i+1} = (1 + \xi)/2$
- $x = x_i \phi_i + x_{i+1} \phi_{i+1}$ is correct at nodes
- Bilinear functions for 2D element have the form $(1 \pm \xi)(1 \pm \eta)/2$

$$\xi = \frac{x - \frac{x_1 + x_2}{2}}{\frac{x_2 - x_1}{2}}$$

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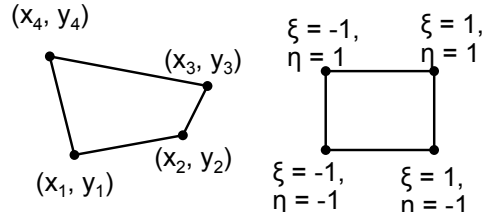
Two-dimensional Element



- Use dimensionless ξ - η coordinate system for shape functions
- Simplest element has one shape function, ϕ_i , at each corner

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Bilinear Basis Functions



Note:
 $\phi_i(\mathbf{x}_{(j)}) = \delta_{ij}$

$$\varphi_1 = \frac{(1-\xi)(1-\eta)}{4} \quad \varphi_2 = \frac{(1+\xi)(1-\eta)}{4}$$

$$\varphi_3 = \frac{(1+\xi)(1+\eta)}{4} \quad \varphi_4 = \frac{(1-\xi)(1+\eta)}{4}$$

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Using Shape Functions

- Model geometry and dependent variable over the element
- Use of same shape functions for both is called isoparametric element
- Shape functions associated with element nodes such that $\phi_i(\mathbf{x}_{(j)}) = \delta_{ij}$

$$x = \sum_{i=1}^4 x_i \varphi_i \quad y = \sum_{i=1}^4 y_i \varphi_i \quad \hat{u} = \sum_{i=1}^4 u_i \varphi_i$$

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Modeling Differential Equation

- Look at simple one-dimensional example: $d^2u/dx^2 + a^2u = 0$
- Equation for u in terms of shape functions gives approximate value $\hat{u} = \sum_{i=1}^N u_i \varphi_i$
- Seek solution in which differential equation is satisfied in an average way over the region; w_i is weighting function

$$\int_0^L w_i(x) \left[\frac{d^2 \hat{u}}{dx^2} + a^2 \hat{u} \right] dx = 0 \quad i = 0, \dots, N$$

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Modeling Differential Equation II

- Last equation on previous slide is general MWR equation for one dimension
- Various choices are used for weighting functions, w_i
- Galerkin method uses $w_i = \phi_i$
- Known to match variational results for certain linear problems

$$\int_0^L \varphi_i(x) \left[\frac{d^2 \hat{u}}{dx^2} + a^2 \hat{u} \right] dx = 0 \quad i = 0, \dots, N$$

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Modeling Differential Equation III

- Use integration by parts to eliminate second derivatives which give zero for linear shape functions

$$\int_0^L \varphi_i \frac{d^2 \hat{u}}{dx^2} dx = \int_0^L \varphi_i \frac{d}{dx} \left(\frac{d \hat{u}}{dx} \right) dx = \int_0^L \varphi_i' \left(\frac{d \hat{u}}{dx} \right) dx$$

$$= \left[\varphi_i \frac{d \hat{u}}{dx} \right]_0^L - \int_0^L \frac{d \varphi_i}{dx} \frac{d \hat{u}}{dx} dx = \left[\varphi_i \frac{d \hat{u}}{dx} \right]_0^L - \int_0^L \frac{d \hat{u}}{dx} \frac{d \varphi_i}{dx} dx$$

- Remember there are N+1 of these equations where $i = 0, \dots, N$

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Modeling Differential Equation IV

- With result of integration by parts, the original equation is (for each $i = 0, \dots, N$)

$$\int_0^L \varphi_i \left[\frac{d^2 \hat{u}}{dx^2} + a^2 \hat{u} \right] dx = \left[\varphi_i \frac{d\hat{u}}{dx} \right]_0^L - \int_0^L \frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} dx$$

$$+ \int_0^L \varphi_i a^2 \hat{u} dx = \left[\varphi_i \frac{d\hat{u}}{dx} \right]_0^L - \int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \hat{u} \right] dx = 0$$

- Next steps: handle boundary terms, separate into elements, and introduce shape functions for u

Modeling Differential Equation V

- At $x = 0$, $\phi_0 = 1$ and all other $\phi_i = 0$
- At $x = L$, $\phi_N = 1$ and all other $\phi_i = 0$
- Have special equations for $i = 0$ and $i = N$

$$i = 0 \quad - \left(1\right) \frac{d\hat{u}}{dx} \Big|_{x=0} - \int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_0}{dx} - \varphi_0 a^2 \hat{u} \right] dx = 0$$

$$i = N \quad \left(1\right) \frac{d\hat{u}}{dx} \Big|_{x=L} - \int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_N}{dx} - \varphi_N a^2 \hat{u} \right] dx = 0$$

$$\text{All other } i \quad \int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \hat{u} \right] dx = 0$$

Modeling Differential Equation VI

$$\int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \hat{u} \right] dx = \int_0^L \frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} dx - \int_0^L \varphi_i a^2 \hat{u} dx = 0$$

- Substitute shape function equation for u into integrals $\hat{u} = \sum_{j=0}^N u_j \varphi_j$

$$\int_0^L \frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} dx = \int_0^L \frac{d \sum_{j=0}^N u_j \varphi_j}{dx} \frac{d\varphi_i}{dx} dx = \sum_{j=0}^N u_j \int_0^L \frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} dx$$

$$\int_0^L \varphi_i a^2 \hat{u} dx = \int_0^L \left[\varphi_i a^2 \sum_{j=0}^N u_j \varphi_j \right] dx = \sum_{j=0}^N u_j \int_0^L \varphi_i a^2 \varphi_j dx$$

Modeling Differential Equation VII

$$\int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \hat{u} \right] dx = \int_0^L \frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} dx - \int_0^L \varphi_i a^2 \hat{u} dx = 0$$

- Result of previous integration is

$$\int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \hat{u} \right] dx = \sum_{j=0}^N u_j \int_0^L \left[\frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \varphi_j \right] dx = 0$$

- This is a system of simultaneous linear equations in the unknowns u_j

$$\sum_{j=0}^N A_{ij} u_j = \sum_{j=0}^N u_j \int_0^L \left[\frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \varphi_j \right] dx = 0$$

Modeling Summary

- For the ODE $u'' + a^2 u = 0$
 - Get a grid of nodes, x_i
 - Each node has a shape function, $\phi_i(x)$ or $N_i(x)$ with the property that $f_i(x_j) = \delta_{ij}$
 - We want to find the value u_i at each node
 - The algebraic equations to do this are

$$\sum_{j=0}^N u_j = 0 \quad \text{where} \quad A_{ij} = \int_0^L \left[\frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \varphi_j \right] dx$$

- Choose ϕ_i and evaluate A_{ij}
 - Linear equations for ϕ_i are easiest

Linear Shape (Basis) Functions

$$\varphi_i(x) = \begin{cases} 0 & x \leq x_{i-1} \\ \frac{x - x_{i-1}}{x_i - x_{i-1}} & x_{i-1} \leq x \leq x_i \\ \frac{x_{i+1} - x}{x_{i+1} - x_i} & x_i \leq x \leq x_{i+1} \\ 0 & x \geq x_{i+1} \end{cases}$$

Get needed shape functions by substituting $i-1$ and $i+1$ for i

Substitute shape functions and derivatives into integrals

Introduction to Elements

- Note that all shape functions are defined for entire region, but are zero for all but a small part of the region
 - Most A_{ij} coefficients will be zero
- In the 1D case we have integral from 0 to L, but basis function, f_j , is nonzero only between x_{i-1} and x_{i+1}
- Look at element between x_i and x_{i+1} where there are only two nonzero element basis functions

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Element Analysis

- Same analysis for individual elements

$$\sum_{j=0}^N A_{ij} u_j \quad A_{ij} = \int_0^L \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx$$

- Look at one typical element where there are only two nonzero shape functions

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Element Shape Functions

- Use element numbering scheme

$$\sum_{j=I}^{II} A_{ij} u_j = 0 \quad A_{ij} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx \quad i = I, II$$

$$\phi_I(x) = \frac{x_{II} - x}{x_{II} - x_I}$$

$$\phi_{II}(x) = \frac{x - x_I}{x_{II} - x_I}$$

- Have two equations for element: $A_{I,I} u_I + A_{I,II} u_{II} = 0$ and $A_{II,I} u_I + A_{II,II} u_{II} = 0$

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One Element Integral

$$\sum_{j=I}^{II} A_{ij} u_j = \sum_{j=I}^{II} u_j \int_{x_I}^{x_{II}} \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx = 0 \quad i = I, II$$

$$A_{ij} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx \quad \phi_I(x) = \frac{x_{II} - x}{x_{II} - x_I}$$

$$A_{I,II} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_{II}}{dx} \frac{d\phi_I}{dx} - \phi_I a^2 \phi_{II} \right] dx \quad \phi_{II}(x) = \frac{x - x_I}{x_{II} - x_I}$$

$$A_{I,I} = \int_{x_I}^{x_{II}} \left[\frac{1}{x_{II} - x_I} \left(\frac{-1}{x_{II} - x_I} \right) - \frac{x_{II} - x}{x_{II} - x_I} a^2 \frac{x - x_I}{x_{II} - x_I} \right] dx$$

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Element Integral Results

- $A_{m,n}$ coefficients for element use α and β notation to summarize results

$$A_{I,I} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_I}{dx} \frac{d\phi_I}{dx} - \phi_I a^2 \phi_I \right] dx = \alpha_I \quad \alpha_I = -\frac{1}{x_{II} - x_I}$$

$$A_{I,II} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_{II}}{dx} \frac{d\phi_I}{dx} - \phi_I a^2 \phi_{II} \right] dx = \beta_I \quad \beta_I = \frac{a^2}{3} (x_{II} - x_I)$$

$$A_{II,I} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_I}{dx} \frac{d\phi_{II}}{dx} - \phi_{II} a^2 \phi_I \right] dx = \beta_{II} \quad \beta_{II} = \frac{1}{x_{II} - x_I}$$

$$A_{II,II} = \int_{x_I}^{x_{II}} \left[\frac{d\phi_{II}}{dx} \frac{d\phi_{II}}{dx} - \phi_{II} a^2 \phi_{II} \right] dx = \alpha_{II} \quad \alpha_{II} = \frac{a^2}{6} (x_{II} - x_I)$$

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Element Assembly

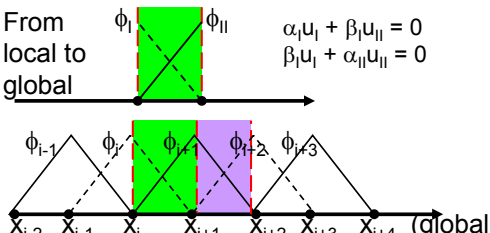
- Return to global numbering scheme
- Two equations for each element
 - $\alpha_I u_I + \beta_I u_{II} = 0$ and $\beta_{II} u_I + \alpha_{II} u_{II} = 0$
- Weight function $i + 1$ produces two element equations that are really one global equation
 - Node I in element from $i + 1$ to $i + 2$
 - Node II in element from i to $i + 1$

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Element Assembly II

- From local to global



$\alpha_i u_i + \beta_i u_{i+1} = 0$
 $\beta_i u_i + \alpha_{i+1} u_{i+1} = 0$
- Weight function ϕ_{i+1} is ϕ_{i+1} in element from x_i to x_{i+1} and ϕ_i from x_{i+1} to x_{i+2}
 - Combine both for ϕ_{i+1} equation

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Element Assembly III

- Combine element equations from some global weight function at each node
- Use node $i + 1$ as general global node
 - $\alpha_i u_i + \beta_i u_{i+1} = 0$ from element $i + 1$ to $i + 2$ becomes $\alpha_{i+1} u_{i+1} + \beta_{i+1} u_{i+2} = 0$
 - $\beta_i u_i + \alpha_{i+1} u_{i+1} = 0$ from element i to $i + 1$ becomes $\beta_i u_i + \alpha_i u_{i+1} = 0$
 - Adding both equations gives $\beta_i u_i + (\alpha_i + \alpha_{i+1}) u_{i+1} + \beta_{i+1} u_{i+2} = 0$
- In 1D problems, can also use global approach (see supplementary charts)

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Equations to be Solved

$$\alpha_0 u_0 + \beta_0 u_1 = - \frac{du}{dx} \Big|_{x=x_0} \quad \alpha_{N-1} u_{N-1} + \beta_{N-1} u_N = \frac{du}{dx} \Big|_{x=x_N}$$

$$\beta_i u_{i-1} + (\alpha_i + \alpha_{i+1}) u_i + \beta_{i+1} u_{i+1} = 0 \quad i = 1, \dots, N-1$$

- Tridiagonal system of $N+1$ equations with $N+3$ variables
 - $N+1$ temperature values and 2 boundary gradients
 - Boundary conditions will specify two other equations

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Boundary Gradients

- If we have Dirichlet boundary conditions, we can solve for temperatures then find gradients
- For Neumann or mixed boundary conditions, we must include gradients in tridiagonal solution
- Write boundary conditions as a $du/dx + b u = c$ and make $g_0 = du/dx|_{x=0}$ the first variable and $g_L = du/dx|_{x=L}$ the last one

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Finite Element Equations

- Equations below only handle boundary conditions with gradients ($a_0 \neq 0$ and $a_N \neq 0$)

$$\begin{bmatrix} a_0 & b_0 & 0 & 0 & 0 & \dots & \dots & 0 & 0 \\ 1 & \alpha_0 + \alpha_1 & \beta_1 & 0 & 0 & \dots & \dots & 0 & 0 \\ & \beta_1 & \alpha_1 + \alpha_2 & \beta_2 & 0 & \dots & \dots & 0 & 0 \\ 0 & 0 & \beta_2 & \alpha_2 + \alpha_3 & \beta_3 & \dots & \dots & 0 & 0 \\ 0 & 0 & 0 & \beta_3 & \ddots & \ddots & \ddots & \vdots & \vdots \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & \dots & \beta_{N-1} & 0 & u_{N-1} \\ 0 & 0 & 0 & 0 & \dots & \dots & \beta_{N-1} & \alpha_{N-1} + \alpha_N & -1 \\ 0 & 0 & 0 & 0 & \dots & \dots & 0 & b_N & a_N \end{bmatrix} \begin{bmatrix} g_0 \\ u_0 \\ u_1 \\ u_2 \\ \vdots \\ \vdots \\ u_{N-1} \\ u_N \\ g_N \end{bmatrix} = \begin{bmatrix} c_0 \\ 0 \\ 0 \\ 0 \\ \vdots \\ \vdots \\ 0 \\ 0 \\ c_N \end{bmatrix}$$

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Solution Errors for $a = 2$

| N | 100 | 100 | 10 | 10 |
|---------------|----------------------|----------------------|----------------------|----------------------|
| Method | FD | FE | FD | FE |
| e_{RMS} | 1.7×10^{-5} | 1.7×10^{-5} | 1.8×10^{-3} | 1.8×10^{-3} |
| e_{max} | 2.4×10^{-5} | 2.4×10^{-5} | 2.4×10^{-3} | 2.4×10^{-3} |
| $e_{grad}(0)$ | 3.6×10^{-4} | 7.0×10^{-5} | 3.6×10^{-2} | 7.0×10^{-3} |
| $e_{grad}(L)$ | 2.1×10^{-4} | 9.6×10^{-5} | 1.8×10^{-2} | 9.5×10^{-3} |

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Solution Errors for a = 0.2

| | | | | |
|----------------------|-----------------------|-----------------------|----------------------|----------------------|
| N | 100 | 100 | 10 | 10 |
| Method | FD | FE | FD | FE |
| e _{RMS} | 6.2x10 ⁻¹⁰ | 6.2x10 ⁻¹⁰ | 6.5x10 ⁻⁸ | 6.5x10 ⁻⁸ |
| e _{max} | 8.6x10 ⁻¹⁰ | 8.6x10 ⁻¹⁰ | 8.5x10 ⁻⁸ | 8.5x10 ⁻⁸ |
| e _{grad(0)} | 1.3x10 ⁻⁶ | 2.2x10 ⁻⁹ | 1.3x10 ⁻⁴ | 2.2x10 ⁻⁷ |
| e _{grad(L)} | 1.3x10 ⁻⁶ | 4.5x10 ⁻⁹ | 1.3x10 ⁻⁴ | 4.4x10 ⁻⁷ |

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- ### Notes on the Error
- The formulations used here for finite elements and finite differences have second order error
 - Notes both equations almost the same
 - Although temperature errors are similar, finite elements gives smaller errors in the gradients
 - The heat source parameter, $a^2 = b/k$, can change the error for a given h
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- ### Higher Order Shape Functions
- Both finite differences and finite elements can decrease grid size to improve accuracy
 - Finite differences can use $O(h)$, $O(h^2)$, $O(h^3)$, etc. expressions for more accuracy with same grid spacing
 - Finite elements can use higher-order shape functions to give more accuracy with same element size
 - Details on supplemental charts
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Quadratic Solution Results

- Solve $d^2u/dx^2 + a^2u = 0$ with Dirichlet boundary conditions with $a = 2$
- 4th order error by log error vs. log h plot

| Number of Elements | Element width, h | Maximum Error | RMS Error |
|--------------------|------------------|---------------|-----------|
| 5 | 0.2 | 3.70E-05 | 2.40E-05 |
| 10 | 0.1 | 2.40E-06 | 1.50E-06 |
| 20 | 0.05 | 1.50E-07 | 9.10E-08 |
| 40 | 0.025 | 9.50E-09 | 5.90E-09 |

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- ### Supplemental Charts
- Charts 48 to 52 show use of global approach to one-dimensional problem with linear shape functions
 - Gives same results as element analysis
 - Demonstrates need for assembly process
 - Not readily accomplished in 2D and 3D
 - Charts 53 to 63 show details of analysis with quadratic shape functions
 - Same overall approach to getting $\sum A_{ij}u_j = 0$, but different equations for ϕ_i
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Global Approach

- Get basis function for each i

$$\int_0^L \left[\frac{d\hat{u}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \hat{u} \right] dx = \sum_{j=0}^N u_j \int_0^L \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx = 0$$

Linear basis functions

- Evaluate integral for each shape function, ϕ_i , to get an equation for each unknown, u_j

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Global Interals

- Value of ϕ_i is nonzero only for $x_{i-1} \leq x \leq x_{i+1}$
- u_j coefficient zero unless $j = i-1, i,$ or $i+1$

$$\sum_{j=0}^N u_j \int_0^L \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx = A_{i,i-1} u_{i-1} + A_{i,i} u_i + A_{i,i+1} u_{i+1} = 0$$

$$A_{ii-1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i-1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i-1} \right] dx$$

$$A_{ii} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_i}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_i \right] dx$$

$$A_{ii+1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i+1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i+1} \right] dx$$

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Results of Integration

$$\alpha_i = \frac{a^2}{3} (x_{i+1} - x_i) - \frac{1}{x_{i+1} - x_i} \quad \beta_i = \frac{a^2}{6} (x_{i+1} - x_i) + \frac{1}{x_{i+1} - x_i}$$

$$\sum_{j=0}^N u_j \int_0^L \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx = A_{i,i-1} u_{i-1} + A_{i,i} u_i + A_{i,i+1} u_{i+1} = 0$$

$$A_{i,i-1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i-1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i-1} \right] dx = \beta_{i-1}$$

$$A_{i,i} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_i}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_i \right] dx = \alpha_i + \alpha_{i-1}$$

$$A_{i,i+1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i+1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i+1} \right] dx = \beta_i$$

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Evaluating the Integrals

- Use A_{ii-1} as an example

$$A_{ii-1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i-1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i-1} \right] dx$$

$$A_{ii-1} = \int_{x_{i-1}}^{x_i} \left[\frac{-1}{x_i - x_{i-1}} \frac{1}{x_i - x_{i-1}} - \frac{x - x_{i-1}}{x_i - x_{i-1}} a^2 \frac{x_i - x}{x_i - x_{i-1}} \right] dx$$

$$\phi_i(x) = \begin{cases} \frac{x - x_{i-1}}{x_i - x_{i-1}} & x_{i-1} \leq x \leq x_i \\ \frac{x_{i+1} - x}{x_{i+1} - x_i} & x_i \leq x \leq x_{i+1} \end{cases} \quad \phi_{i-1}(x) = \begin{cases} \frac{x - x_{i-2}}{x_{i-1} - x_{i-2}} & x_{i-2} \leq x \leq x_{i-1} \\ \frac{x_i - x}{x_i - x_{i-1}} & x_{i-1} \leq x \leq x_i \end{cases}$$

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Constant Steps $x_{i+1} - x_i = h$

$$\alpha_i = \frac{ha^2}{3} - \frac{1}{h} \quad \beta_i = \frac{ha^2}{6} + \frac{1}{h}$$

$$\sum_{j=0}^N u_j \int_0^L \left[\frac{d\phi_j}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_j \right] dx = A_{i,i-1} u_{i-1} + A_{i,i} u_i + A_{i,i+1} u_{i+1} = 0$$

$$A_{i,i-1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i-1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i-1} \right] dx = \frac{ha^2}{6} + \frac{1}{h}$$

$$A_{i,i} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_i}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_i \right] dx = \frac{2ha^2}{3} - \frac{2}{h}$$

$$A_{i,i+1} = \int_{x_{i-1}}^{x_{i+1}} \left[\frac{d\phi_{i+1}}{dx} \frac{d\phi_i}{dx} - \phi_i a^2 \phi_{i+1} \right] dx = \frac{ha^2}{6} + \frac{1}{h}$$

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Quadratic Shape Functions

- Use one-dimensional shape functions here as an example
- Need three nodes per element for quadratic shape function
 - Label nodes as $i - 1, i,$ and $i + 1$
 - Define coordinates as shown below

$$h_i = x_i - x_{i-1} = x_{i+1} - x_i = \frac{x_{i+1} - x_{i-1}}{2}$$

$$\xi = 2 \frac{x - x_i}{x_{i+1} - x_{i-1}} \quad \xi = \frac{x - x_i}{h_i}$$

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Quadratic Basis Functions II

$$\phi_{i+1}(\xi) = \frac{\xi(\xi+1)}{2} = \frac{1}{2} \frac{x - x_i}{h_i} \left(\frac{x - x_i}{h_i} + 1 \right)$$

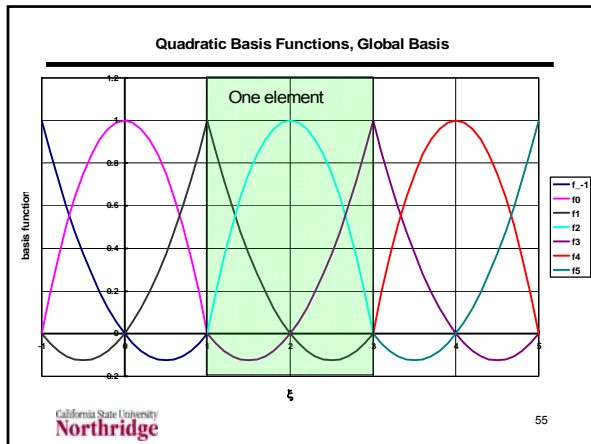
$$\phi_i(\xi) = 1 - \xi^2 = 1 - \left(\frac{x - x_i}{h_i} \right)^2$$

$$\phi_{i-1}(\xi) = \frac{\xi(\xi-1)}{2} = \frac{1}{2} \frac{x - x_i}{h_i} \left(\frac{x - x_i}{h_i} - 1 \right)$$

$$h_i = x_i - x_{i-1} = x_{i+1} - x_i = \frac{x_{i+1} - x_{i-1}}{2}$$

- Can use either definition of h_i in integrals
- Pick one to give simplest results

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Quadratics for $d^2u/dx^2+a^2u = 0$

- Have three u_j coefficients in each element

$$\sum_{j=I}^{III} u_j \int_{x_i}^{x_{i+1}} \left[\frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \varphi_j \right] dx = 0 \quad i = I, II, III$$

- Have three shape functions in each element
- Nine integrals to evaluate (3 u_j coefficients times 3 basis functions)
 - Three found by symmetry

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Quadratic Element Results

- Result of element integrals

$$\sum_{j=I}^{III} u_j \int_{x_i}^{x_{i+1}} \left[\frac{d\varphi_j}{dx} \frac{d\varphi_i}{dx} - \varphi_i a^2 \varphi_j \right] dx = 0 \quad i = I, II, III$$

$$\begin{bmatrix} \frac{7}{6h_i} - \frac{4a^2h_i}{15} & \frac{4}{3h_i} - \frac{2a^2h_i}{15} & \frac{1}{6h_i} + \frac{a^2h_i}{15} \\ \frac{4}{3h_i} - \frac{2a^2h_i}{15} & \frac{7}{6h_i} - \frac{4a^2h_i}{15} & \frac{4}{3h_i} - \frac{2a^2h_i}{15} \\ \frac{1}{6h_i} + \frac{a^2h_i}{15} & \frac{4}{3h_i} - \frac{2a^2h_i}{15} & \frac{7}{6h_i} - \frac{4a^2h_i}{15} \end{bmatrix} \begin{bmatrix} u_I \\ u_{II} \\ u_{III} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

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Quadratic Assembly

- Nodes from $i = 0$ to $i = N$
- Elements from $k = 0$ to $K = 2N - 2$
- Element k has nodes $2k, 2k+1$ and $2k+2$
- Different equations for odd and even i

$$\left[-\frac{4}{3h_k} - \frac{2a^2h_k}{15} \right] u_{i-1} + \left[\frac{8}{3h_k} - \frac{16a^2h_k}{15} \right] u_i + \left[-\frac{4}{3h_k} - \frac{2a^2h_k}{15} \right] u_{i+1} = 0 \quad \text{Odd } i, k = \frac{i-1}{2}$$

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Quadratic Assembly II

- Even i equation

$$\left[\frac{1}{6h_{k-1}} + \frac{a^2h_{k-1}}{15} \right] u_{i-2} + \left[-\frac{4}{3h_{k-1}} - \frac{2a^2h_{k-1}}{15} \right] u_{i-1} + \left[\frac{7}{6h_{k-1}} - \frac{4a^2h_{k-1}}{15} + \frac{7}{6h_k} - \frac{4a^2h_k}{15} \right] u_i + \left[-\frac{4}{3h_k} - \frac{2a^2h_k}{15} \right] u_{i+1} + \left[\frac{1}{6h_k} + \frac{a^2h_k}{15} \right] u_{i+2} = 0$$

Even i
 $k = \frac{i}{2}$

- Boundary nodes on next chart

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Quadratic Assembly III

- Boundary nodes at $i = 0$ and $i = N$

$$\left[\frac{7}{6h_0} - \frac{4a^2h_0}{15} \right] u_0 + \left[-\frac{4}{3h_0} - \frac{2a^2h_0}{15} \right] u_1 + \left[\frac{1}{6h_0} + \frac{a^2h_0}{15} \right] u_2 = -\frac{du}{dx} \Big|_{x=0}$$

$$\left[\frac{7}{6h_k} - \frac{4a^2h_k}{15} \right] u_N + \left[-\frac{4}{3h_k} - \frac{2a^2h_k}{15} \right] u_{N-1} + \left[\frac{1}{6h_k} + \frac{a^2h_k}{15} \right] u_{N-2} = \frac{du}{dx} \Big|_{x=L}$$

- Can simplify equations for constant element size, h_k

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Constant h

- Multiply all results by 3h or 6h to get whole number in lead term and a²h² in remaining terms
 - Even i result below with odd i, i = 0, and i = N results on next chart

$$\left[1 + \frac{6a^2h^2}{15}\right]u_{i-2} + \left[-8 - \frac{12a^2h^2}{15}\right]u_{i-1} + \left[14 - \frac{48a^2h^2}{15}\right]u_i + \left[-8 - \frac{12a^2h^2}{15}\right]u_{i+1} + \left[1 + \frac{6a^2h^2}{15}\right]u_{i+2} = 0 \quad \text{Even } i$$

Constant h (i = 0, N and odd i)

$$\left[7 - \frac{24a^2h^2}{15}\right]u_0 + \left[-8 - \frac{12a^2h^2}{15}\right]u_1 + \left[1 + \frac{6a^2h^2}{15}\right]u_2 = -h \frac{du}{dx} \Big|_{x=0}$$

$$\left[-4 - \frac{6a^2h^2}{15}\right]u_{i-1} + \left[8 - \frac{48a^2h^2}{15}\right]u_i + \left[-4 - \frac{6a^2h^2}{15}\right]u_{i+1} = 0 \quad \text{Odd } i$$

$$\left[7 - \frac{4a^2h^2}{15}\right]u_N + \left[-8 - \frac{2a^2h^2}{15}\right]u_{N-1} + \left[1 + \frac{a^2h^2}{15}\right]u_{N-2} = h \frac{du}{dx} \Big|_{x=L}$$

Quadratic Solutions

- Have N + 1 equations in N + 3 unknowns
- Add two boundary conditions equations
- Set of equations forms pentadiagonal matrix
- Solve by extension of Thomas algorithm that applies Gauss elimination to matrix with five diagonals next to each other