

Numerical Solutions of Elliptic Equations

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Seminar in Engineering Analysis

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Outline

- Review last class
- Numerical solution of elliptic equations
- Laplace's equation as an example
- Finite difference forms
- Iterative solution of algebraic equations
- Sample results

Review Properties of Solutions

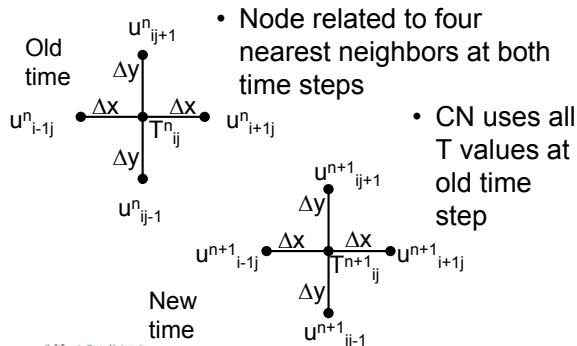
- *Consistency* – truncation error becomes zero as step sizes approach zero
- *Stability* – errors remain bounded
- *Convergent* – tends to the exact solution of the differential equation as the grid size tends towards zero
- *Physical reality* – Solutions produce physically realistic results
- *Accuracy* – Many sources of error in numerical solutions.

Review Two Space Dimensions

$$\left[\frac{\partial T}{\partial t} = \alpha \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) \right]_{ij}^{(time\ step=n,n+1/2,n+1)}$$

- Extension of one space dimension
- Have grid in y_j as well as x_i and time
- Explicit method has almost no difference. but has more stringent stability limit
- Cannot apply Thomas algorithm directly to implicit method
- Define $f_x = \alpha \Delta t / (\Delta x)^2$ and $f_y = \alpha \Delta t / (\Delta y)^2$

Finite-difference Grid



Review Algorithms

- Explicit method $T_{ij}^{n+1} = f_x (T_{i+1j}^n + T_{i-1j}^n) + f_y (T_{ij+1}^n + T_{ij-1}^n) + (1 - 2f_x - 2f_y) T_{ij}^n$
- Crank-Nicholson $-f_y T_{ij+1}^{n+1} - f_x T_{i+1j}^{n+1} + 2(1 + f_x + f_y) T_{ij}^{n+1} - f_x T_{i-1j}^{n+1} - f_y T_{ij-1}^{n+1} = f_x (T_{i+1j}^n + T_{i-1j}^n) + f_y (T_{ij+1}^n + T_{ij-1}^n) + 2(1 - f_x - f_y) T_{ij}^n = R_{ij}^n$
- Fully implicit $-f_y T_{ij-1}^{n+1} - f_x T_{i-1j}^{n+1} + (1 + 2f_x + 2f_y) T_{ij}^{n+1} - f_x T_{i+1j}^{n+1} - f_y T_{ij+1}^{n+1} = T_{ij}^n$

Review Explicit Stability

- $1 - 2f_x - 2f_y$ must be positive for stability so $f_x + f_y$ must be less than 1/2

$$T_{ij}^{n+1} = f_x(T_{i+1j}^n + T_{i-1j}^n) + f_y(T_{ij+1}^n + T_{ij-1}^n) + (1 - 2f_x - 2f_y)T_{ij}^n$$

$$f_x + f_y = \frac{\alpha\Delta t}{(\Delta x)^2} + \frac{\alpha\Delta t}{(\Delta y)^2} \leq \frac{1}{2}$$

- To cut Δx and Δy by a factor of 2 we must cut Δt by a factor of 4
- Work increases by a factor of 16

Alternating Direction Implicit

- ADI is technique for allowing use of Thomas algorithm for simple solution
 - At each time step, treat one direction as explicit and one as implicit
 - Even number of time steps required

$$\frac{u_{ij}^* - u_{ij}^n}{\Delta t} = \alpha \left[\frac{u_{i+1j}^* + u_{i-1j}^* - 2u_{ij}^*}{(\Delta x)^2} + \frac{u_{ij+1}^n + u_{ij-1}^n - 2u_{ij}^n}{(\Delta y)^2} \right]$$

$$\frac{u_{ij}^{n+2} - u_{ij}^*}{\Delta t} = \alpha \left[\frac{u_{i+1j}^* + u_{i-1j}^* - 2u_{ij}^*}{(\Delta x)^2} + \frac{u_{ij+1}^{n+2} + u_{ij-1}^{n+2} - 2u_{ij}^{n+2}}{(\Delta y)^2} \right]$$

Elliptic Equations

- Solutions at any point in the region depend on conditions at all boundaries
- No open boundaries like parabolic case
- Main examples are Laplace and Poisson Equations

$$\nabla^2 u = 0 \quad \nabla^2 u = F(\text{space})$$

- Laplace equation in two-dimensional Cartesian coordinates $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$

2D Cartesian Laplace

- Use finite-difference grid where $x_i = x_0 + i\Delta x$ and $y_j = y_0 + j\Delta y$
 - $x_0 \leq x_i \leq x_N$ and $y_0 \leq y_j \leq y_M$
 - $0 \leq i \leq N$ and $0 \leq j \leq M$
 - $\Delta x = (x_N - x_0)/N$ and $\Delta y = (y_M - y_0)/M$
 - Must specify boundary conditions on u at all boundaries
 - South $j = 0, 0 \leq i \leq N$
 - North $j = M, 0 \leq i \leq N$
 - West $i = 0, 0 \leq j \leq M$
 - East $i = N, 0 \leq j \leq M$
- Dependent variable, u
 $u_{ij} = u(x_i, y_j)$

Finite Differences

- Use second-order expressions for second derivatives

$$\frac{\partial^2 u}{\partial x^2} \Big|_{ij} = \frac{u_{i+1j} + u_{i-1j} - 2u_{ij}}{(\Delta x)^2} + O[(\Delta x)^2]$$

$$\frac{\partial^2 u}{\partial y^2} \Big|_{ij} = \frac{u_{ij+1} + u_{ij-1} - 2u_{ij}}{(\Delta y)^2} + O[(\Delta y)^2]$$

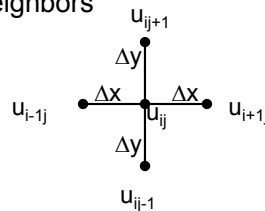
$$\beta = \frac{\Delta x}{\Delta y}$$

$$\left[\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right]_{ij} = \frac{u_{i+1j} + u_{i-1j} - 2u_{ij}}{(\Delta x)^2} + \frac{u_{ij+1} + u_{ij-1} - 2u_{ij}}{(\Delta y)^2} + O[(\Delta x)^2, (\Delta y)^2] = 0$$

- Multiply by $(\Delta x)^2$

Finite-difference Equation

- Links central node to four nearest neighbors



- Only parameter: $\beta = \Delta x/\Delta y$
- For $\beta = 1, 2(1 + \beta^2) = 4$

$$u_{i+1j} + u_{i-1j} + \beta^2(u_{ij+1} + u_{ij-1}) - 2(1 + \beta^2)u_{ij} = 0$$

Finite-difference Equation

- For 2D Cartesian Laplace equation with $\beta = \Delta x/\Delta y = 1$, u_{ij} is the average of its four nearest neighbors

$$u_{ij} = \frac{u_{i+1j} + u_{i-1j} + u_{ij+1} + u_{ij-1}}{4}$$

- Consider Dirichlet boundary conditions
 - u_{ij} known at all boundary nodes
 - Need to find $(N - 1)(M - 1)$ unknown values of u_{ij} at central nodes

General Equation

- Most general case has five diagonals, but can have different terms $\frac{\partial}{\partial x} k \frac{\partial T}{\partial x} + \frac{\partial}{\partial y} k \frac{\partial T}{\partial y} + \dot{Q} = 0$
 - Occurs with variable properties
- Notation A_{ij}^{point} is general coefficient
 - ij refers to a particular node
 - Point = N(orth), S(outh), E(ast), W(est) refers to neighboring nodes by direction
- General equation shown below

$$A_{ij}^S u_{ij-1} + A_{ij}^W u_{i-1j} + A_{ij}^P u_{ij} + A_{ij}^E u_{i+1j} + A_{ij}^N u_{ij+1} = b_{ij}$$

Solving the Equations

- Typically have large number of equations forming sparse matrix
 - For $\Delta x = \Delta y = .01$ have 99^2 equations so matrix has 96×10^6 potential coefficients
 - Only 48609 (0.051%) are nonzero
- Want data structure and algorithm for handling sparse matrices
- Gauss elimination uses storage for banded matrices
- Iterative methods used for solutions

Iterative Solutions

- Simplest examples are Jacobi, Gauss-Seidel, and Successive Over Relaxation
- Move from iteration n to iteration $n+1$
- Iteration 0 is initial guess (often all zero)
- Solve equation for u_{ij} and use this as basis for iteration

$$u_{ij} = \frac{b_{ij}}{A_{ij}^P} - \frac{A_{ij}^S u_{ij-1} + A_{ij}^W u_{i-1j} + A_{ij}^E u_{i+1j} + A_{ij}^N u_{ij+1}}{A_{ij}^P}$$

$$u_{ij} = b'_{ij} - A_{ij}^S u_{ij-1} - A_{ij}^W u_{i-1j} - A_{ij}^E u_{i+1j} - A_{ij}^N u_{ij+1}$$

Iterative Solutions II

- Use superscript (n) for iteration number
- Jacobi iteration uses all old values

$$u_{ij}^{(n+1)} = b'_{ij} - A_{ij}^S u_{ij-1}^{(n)} - A_{ij}^W u_{i-1j}^{(n)} - A_{ij}^E u_{i+1j}^{(n)} - A_{ij}^N u_{ij+1}^{(n)}$$

- Gauss-Seidel uses most-recent values

$$u_{ij}^{(n+1)} = b'_{ij} - A_{ij}^S u_{ij-1}^{(n+1)} - A_{ij}^W u_{i-1j}^{(n+1)} - A_{ij}^E u_{i+1j}^{(n)} - A_{ij}^N u_{ij+1}^{(n)}$$

- Relaxation basis: Gauss-Seidel provides a correction that can be adjusted

$$u_{ij}^{(n+1)} = u_{ij}^{(n)} + \omega [u_{ij}^{(n+1,GS)} - u_{ij}^{(n)}]$$

Relaxation Factor

Relaxation Methods

- Relaxation factor, ω , greater than or less than 1 is over- or underrelaxation
 - Underrelaxation procures stability in problems that will not converge
 - Overrelaxation procures speed in well-behaved problems

$$u_{ij}^{(n+1)} = u_{ij}^{(n)} + \omega [u_{ij}^{(n+1,GS)} - u_{ij}^{(n)}] = (1 - \omega) u_{ij}^{(n)} + \omega u_{ij}^{(n+1,GS)}$$

$$= (1 - \omega) u_{ij}^{(n)} - \omega [A_{ij}^S u_{ij-1}^{(n+1)} + A_{ij}^W u_{i-1j}^{(n+1)} + A_{ij}^E u_{i+1j}^{(n)} + A_{ij}^N u_{ij+1}^{(n)} - b'_{ij}]$$

Excel Relaxation Code I

- Gauss-Seidel solution to Poisson equation with constant source term

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\dot{Q}}{k} = 0$$

$$u_{i+1j} + u_{i-1j} + \beta^2(u_{ij+1} + u_{ij-1}) - 2(1 + \beta^2)u_{ij} + \frac{(\Delta x)^2 \dot{Q}}{k} = 0$$

$$u_{ij} = \frac{u_{i+1j} + u_{i-1j} + \beta^2(u_{ij+1} + u_{ij-1}) + \frac{(\Delta x)^2 \dot{Q}}{k}}{2(1 + \beta^2)} = 0$$
- Iterative Excel Formula for Cell F9
 =(E9+G9+betaSqd*(F8+F10)+Source)/Denominator

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Excel Relaxation Code II

- Modify previous solution to include relaxation factor, ω

$$u_{ij}^{(n+1)} = (1 - \omega)u_{ij}^{(n)} - \omega \frac{u_{i+1j}^{(n)} + u_{i-1j}^{(n)} + \beta^2(u_{ij+1}^{(n)} + u_{ij-1}^{(n)}) + \frac{(\Delta x)^2 \dot{Q}}{k}}{2(1 + \beta^2)}$$
- Iterative Excel Formula for Cell F9
 =oneMinusw*F9+omega*(E9+G9+betaSqd*(F8+F10)+Source)/Denominator

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Visual Basic Relaxation Code

```

for iter = 1 to maxIter
maxResid = 0
for i = 1 to N - 1
for j = 1 to N - 1
old = u(i, j)
u(i, j) = (1 - omega) * u(i, j)
- omega * ( AN(i, j) * u(i, j+1)
+ AE(i, j) * u(i+1, j) + AS(i, j)
* u(i, j-1) + AW(i, j) * u(i-1, j)
- b(i, j) )
resid = abs( ( u(i, j) - old ) /
u(i, j) )
if resid > maxResid then
maxResid = resid
end if

```

One set of iterations (omitted on next page)

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Visual Basic Relaxation Code II

```

for iter = 1 to maxIter
maxResid = 0;
for i = 1 to N - 1
for j = 1 to N - 1
' Code in red from previous slide
next j
next i
if maxResid <= errTol exit for
next iter
if maxResid > errTol then
print #1 "Not converged"
else
call doOutput( u, N, M )
end if

```

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Converging Iterations

- Have three different results
 - Correct solution to differential equation
 - Correct solution finite-difference equations
 - Current and previous iteration values
- Iterations should approach correct solution to finite-difference equations
- Since neither correct solution is known, we use norm of error estimates
 - Residual in finite-difference equations
 - Change in iteration value

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Converging Iterations II

- Take infinity norm (maximum absolute value) or L2 norm (RMS error)
 - Former easiest to code

$$\left[\begin{matrix} \text{Relative} \\ \text{Change} \end{matrix} \right]_{ij} = \left| \frac{u_{ij}^{(n+1)} - u_{ij}^{(n)}}{u_{ij}^{(n+1)}} \right|$$

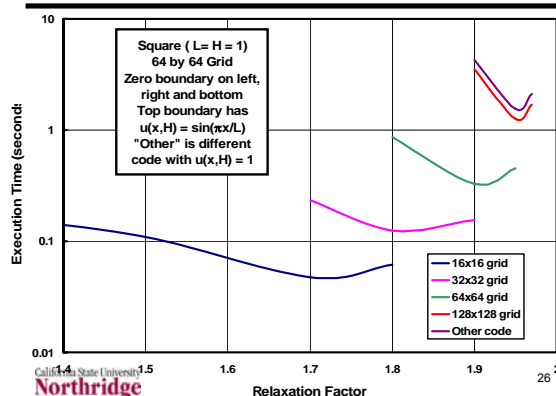
$$\left[\text{Residual} \right]_{ij}^{(n+1)} = \left| u_{ij}^{(n+1)} + A_{ij}^{S'} u_{ij-1}^{(n+1)} + A_{ij}^{W'} u_{i-1j}^{(n+1)} + A_{ij}^{E'} u_{i+1j}^{(n+1)} + A_{ij}^{N'} u_{ij+1}^{(n+1)} - b_{ij}' \right|_{24}$$

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Execution Times and Errors

- Examine square region with zero boundary conditions at $x = 0$, $x = x_{\max}$, and $y = 0$; two cases for $y = y_{\max}$
 - Case 1: constant value of $u_N(x) = 1$
 - Case 2: $u_N(x) = \sin(\pi x)$
- First case has discontinuity for $y = y_{\max}$ at $x = 0$ and $x = x_{\max}$
- Use overrelaxation (SOR) with variable relaxation factors

Effect of Relaxation Factor on Execution Time



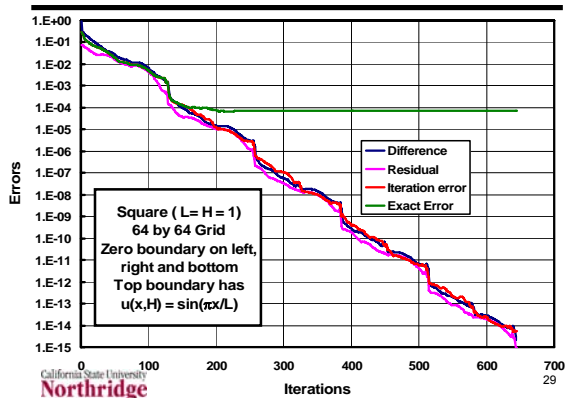
Execution Times and Errors II

- For second order algorithm the error depends on the boundary condition
 - Case 1: constant value of $u_N(x) = 1$
 - Case 2: $u_N(x) = \sin(\pi x)$
- First case has discontinuity for $y = y_{\max}$ at $x = 0$ and $x = x_{\max}$
 - First case gives first order error due to discontinuity
 - Second case gives second order error

Effect of Iterations on Errors

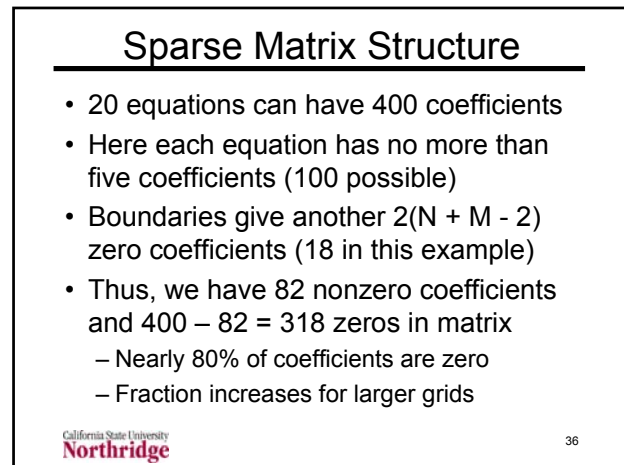
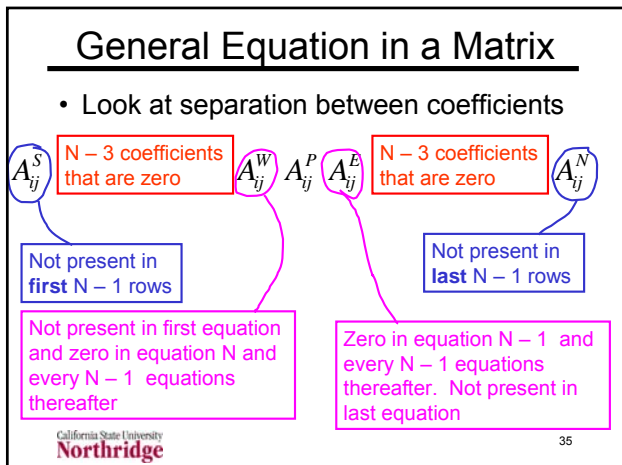
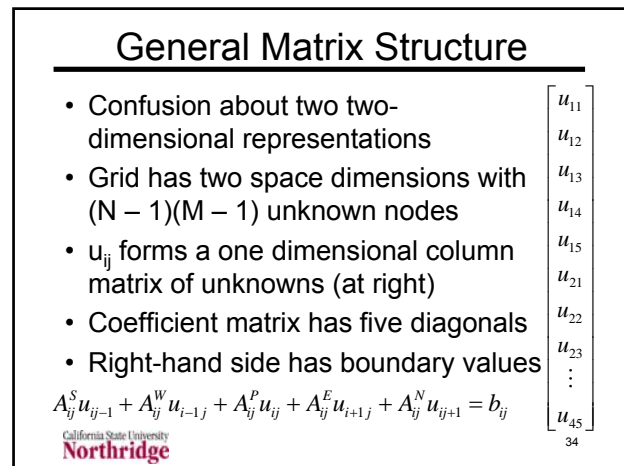
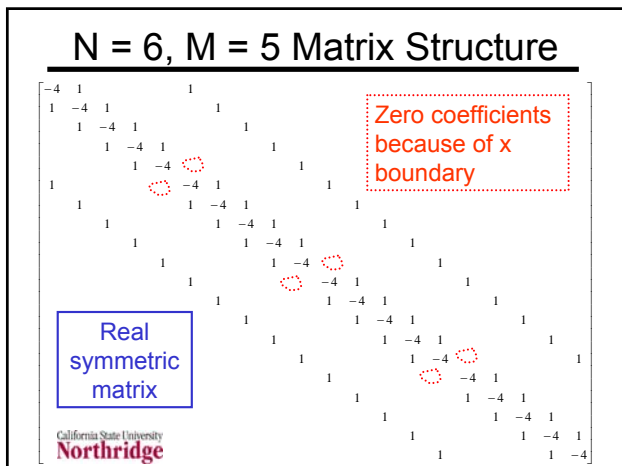
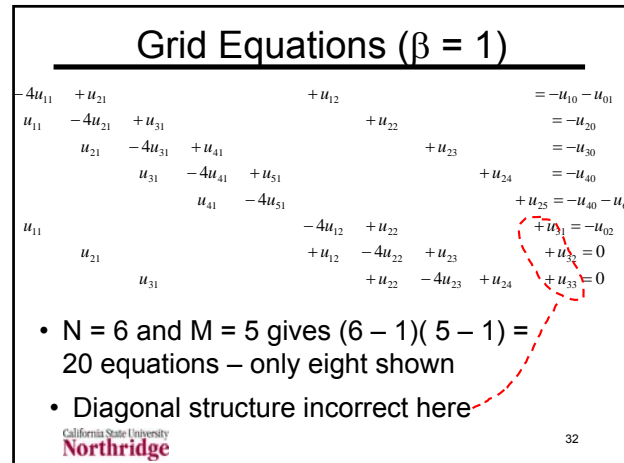
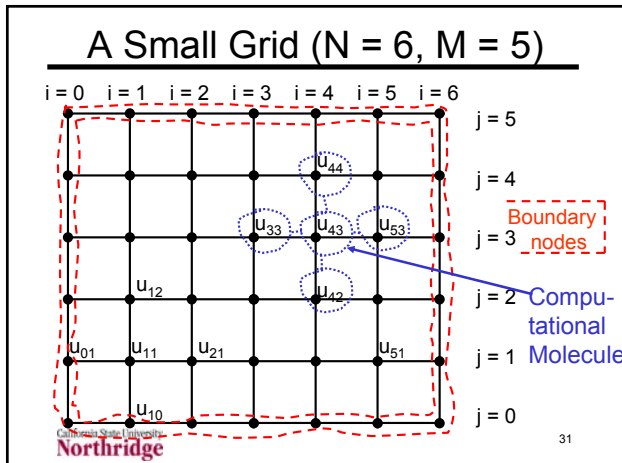
- Compare three error measures using the maximum value on the grid
 - True iteration error: difference between the current value and the value found by an exact solution of the difference equations
 - Difference in u_{ij} between two iterations
 - Residual = $u_{i+1j} + u_{i-1j} + u_{ij+1} + u_{ij-1} - 4u_{ij}$
 - Exact error is difference between iteration value and exact solution

Effects of Iterations on Laplace Equation Errors



Advanced Solvers

- Advanced solution techniques treat matrix for finite-difference equations
- Leads to dimensional confusion
 - Start with 2D grid (x and y indices)
 - Treat as matrix equation where unknowns u_{ij} form a column vector (one-dimensional)
 - The coefficients in the matrix form a two-dimensional display
 - Examine small grid example



How Sparse is the Matrix?

- The M by N grid has $(M - 1)(N - 1)$ nodes with equations giving $(M - 1)^2(N - 1)^2$ possible coefficients
- Without boundaries we have only 5 $(M - 1)(N - 1)$ nonzero coefficients
- Boundaries give $2(N + M - 2) = 2(M - 1) + 2(N - 1)$ additional zero coefficients

$$\left[\begin{array}{l} \text{Nonzero} \\ \text{Fraction} \end{array} \right] = \frac{5(N-1)(M-1) - 2(N-1) - 2(M-1)}{(N-1)^2(M-1)^2} = \frac{5}{(N-1)(M-1)} - \frac{2}{(N-1)(M-1)^2} - \frac{2}{(N-1)^2(M-1)}$$

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What Makes Sparseness?

- Each node is connected only to a small number of nearest neighbors
 - Problem here has four neighbors
 - Higher order schemes and 3D Laplace equation can have more neighbors
- Can have complex coefficients so long as number of neighbors is limited
- Look at uneven grid spacing as an example of this

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