

Unit Seven – Introduction to the Second Law

Mechanical Engineering 370

Thermodynamics

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Outline

- Quiz six and midterm results and solution (Note revised schedule)
- Introduction to the Second Law of Thermodynamics
- Physical background for second law
- Mathematical statement of second law
- Use of entropy as determination of maximum efficiency
- Deriving other common forms of the second law

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Unit Seven Goals

- As a result of studying this unit you should be able to
 - recognize that there is a thermodynamic property, called entropy, s , defined as follows: $ds = (du + Pdv) / T$
 - understand the inequality that $dS \geq dQ/T$
 - understand that $dS \geq 0$ for isolated systems
 - recognize that entropy is a property
 - Understand that the $=$ part of the \geq sign applies in a reversible process

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More Unit Seven Goals

- understand the definitions of engine and refrigeration cycles
- apply the definitions of work and heat flow for these cycles
- compute the efficiency of an engine cycle
- compute the coefficient of performance (COP) for a refrigeration cycle
- perform Carnot cycle computations
- Read second law notes from web site

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Why the Second Law?

- Encapsulates the phenomenon that certain process in nature flow one way
 - Water flows downhill
 - Heat flows from high to low temperatures
- We know that we can reverse these processes with an external effect
 - pump water uphill
 - use a refrigerator to transfer heat from low to high temperature

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Second Law Application

- Based on entropy, a thermodynamic property
- Used to define ideal (reversible) processes
- Shows that maximum efficiency of conversion of heat to work occurs in a reversible process
- Gives quantitative calculations for this maximum efficiency

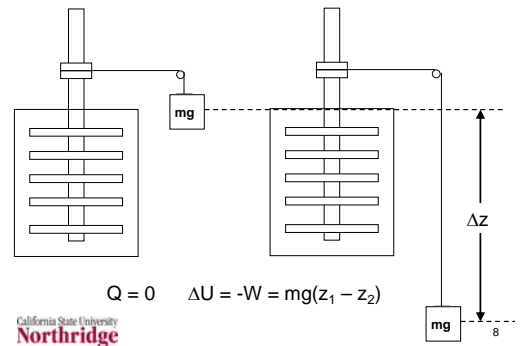
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Approach to Second Law

- Developed around 1850 by considering of engine and refrigeration cycles
- Text considers similar derivation
- Important idea is in result using entropy
- Class notes will start at this point
 - Provide focus on ultimate calculations
 - Show equivalence to text derivation

Joule's Experiment



Joule's Experiment Explained

- The falling weight, mg , turns the paddle wheels increasing the system energy
- For an insulated system $Q = \Delta U + W = 0$ so that $\Delta U = -W = mg(z_{\text{initial}} - z_{\text{final}})$
- For falling weight, $z_{\text{initial}} - z_{\text{final}} > 0$, so $\Delta U > 0$, corresponding to water heating.
- What about process where weight rises and water cools?

Another Example

- Two identical blocks
 - Same mass and heat capacity
 - Block A at 300 K, block B at 500 K
 - Blocks placed in contact each reaching a final temperature of 400 K
 - No heat or work external to blocks
 - $\Delta U = \Delta U_A + \Delta U_B = 0$ or $\Delta U_A = -\Delta U_B$
- Can $T_{A,\text{Final}} = 200 \text{ K}$ and $T_{B,\text{Final}} = 600 \text{ K}$?
 - This also gives $\Delta U = \Delta U_A + \Delta U_B = 0$ or $\Delta U_A = -\Delta U_B$

General Idea

- Some processes in nature are observed to only proceed in one direction
- First law does not prohibit these processes going in the opposite direction
- Is there any general rule that shows the one-directional nature of processes
- Yes, it is the second law

The Second Law

- There exists an extensive thermodynamic property called the entropy, S , defined as follows:

$$dS = (dU + PdV)/T$$
- For any process $dS \geq dQ/T$
- For an isolated system $dS \geq 0$
- T must be absolute temperature
- Inequality is required to show directional effects

Entropy as a Property

- Have total entropy, S , and specific entropy, $s = S/m$
- Dimensions of entropy are energy divided by temperature
- For S , typical units are kJ/K or Btu/R
 - Units for s : kJ/kg·K or Btu/lb_m·R
 - s in tables is similar to v
 - One-phase regions give $s(T,P)$
 - In mixed region, $s = s_f + x s_{fg}$
 - Ideal gas equations for entropy as well

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Cyclic Processes

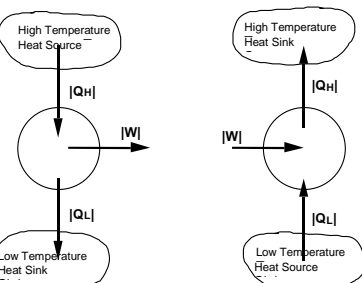
- In a cycle, the initial and final states of the system are the same
- Since the initial and final states are the same, the properties of the initial and final states are the same
- Thus, for a cycle, $\Delta u = \Delta s = 0$
- Since heat and work depend on path, these may be nonzero, but $Q = \Delta U + W$ means that $Q = W$ for a cycle

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Cycles with $|Q_H| = |Q_L| + |W|$

- Engine cycle converts heat to work
- Refrigeration cycle transfers heat from low to high temperature



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Engine Cycle

Refrigeration Cycle

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Engine Cycle Heat and Work

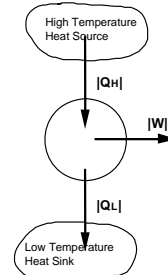
- Relation between cycle terms and usual sign conventions
- Engine cycle

$$|Q_H| = -Q_{HTHS} = Q_{H,Cycle}$$

$$|Q_L| = Q_{LTHS} = -Q_{L,Cycle}$$

$$Q_{cycle} = |Q_H| - |Q_L|$$

$$Q_{cycle} = W > 0$$



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Engine Cycle Schematic

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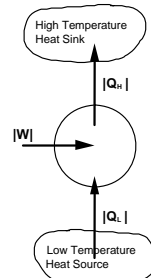
Refrigeration Cycle Heat/ Work

- Relation between cycle terms and usual sign conventions
- Refrigeration cycle

$$|Q_H| = Q_{HTHS} = -Q_{H,Cycle}$$

$$|Q_L| = -Q_{LTHS} = Q_{L,Cycle}$$

$$Q_{cycle} = |Q_L| - |Q_H| = W < 0$$



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Refrigeration Cycle Schematic

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Cycle Parameters

- Engine cycle efficiency $\eta = \frac{|W|}{|Q_H|}$
- Refrigeration cycle COP (coefficient of performance) $COP = \beta = \frac{|Q_L|}{|W|}$
- General definitions, valid for any cycle
- Engine efficiency always less than one
- COP can be greater than one

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Reversible Process

- In a reversible process where $dS_{\text{isolated system}} = 0$ it is possible to return a system to its initial state with no changes in the surroundings
- This is an idealization; we cannot do better than a reversible process
- External reversibility is when $dS = 0$ is true for an isolated system
- Internal reversibility is when $dS = dQ/T$ for one subsystem in an isolated system
 - May still have $dS_{\text{isolated system}} > 0$

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Temperature Reservoir

- A body that transfers heat with no change in its temperature
- Two-phase fluid is best example
- Reservoir usually envisioned as very large body such that $\Delta T = Q/(mc_v) \approx 0$
- Basic idea is that instead of $dS = dQ/T$ we can write $\Delta S = \int dQ/T = Q/T$
- Temperature reservoir is internally reversible

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Carnot Cycle Introduction

- This cycle has theoretical significance but is not a practical cycle
- In traditional approaches to the second law, results for Carnot cycles are used to derive entropy
- Here we will show that Carnot cycle results can be obtained from principle that $dS_{\text{isolated system}} \geq 0$

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Carnot Cycle Results

- Carnot cycle uses one high temperature reservoir at T_H and one low temperature reservoir at T_L
- We have results for Carnot engine cycle efficiency and Carnot refrigeration cycle coefficient of performance
- Efficiency $\eta = |W| / |Q_H| \leq 1 - T_L / T_H = \eta_C$
- COP = $|Q_L| / |W| \leq T_L / (T_H - T_L) = \beta_{\text{Carnot}}$

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Carnot Cycle

- Cycle with only two temperature reservoirs, HTR at T_H and LTR at T_L
- $\Delta S_{\text{isol syst}} = \Delta S_{\text{HTR}} + \Delta S_{\text{cycle}} + \Delta S_{\text{LTR}} \geq 0$
- $\Delta S_{\text{isol syst}} = Q_{\text{HTR}}/T_H + 0 + Q_{\text{LTR}}/T_L \geq 0$
- For the engine cycle $Q_{\text{HTR}} = -|Q_H|$ and $Q_{\text{LTR}} = |Q_L|$, so $-|Q_H|/T_H + |Q_L|/T_L \geq 0$
- Refrigeration cycle $Q_{\text{HTR}} = |Q_H|$ and $Q_{\text{LTR}} = -|Q_L|$, so $|Q_H|/T_H - |Q_L|/T_L \geq 0$

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Carnot Cycle Continued

- For both cycles, $|Q_H| = |Q_L| + |W|$
- For the engine cycle
 - $-|Q_H|/T_H + |Q_L|/T_L \geq 0$ so that $-|Q_H|/T_H + (|Q_H| - |W|)/T_L \geq 0$
 - Rearrange to get $|Q_H|(1/T_L - 1/T_H) \geq |W|/T_L$
 - $\eta = |W| / |Q_H| \leq 1 - T_L / T_H = \eta_{\text{Carnot}}$
- It is impossible to construct an engine cycle without heat rejection

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Carnot Cycle Concluded

- For both cycles, $|Q_H| = |Q_L| + |W|$
- For the refrigeration cycle
 - $|Q_H|/T_H - |Q_L|/T_L \geq 0$ so that $(|Q_L| + |W|)/T_H \geq |Q_L|/T_L$
 - $|Q_L|/T_L \geq 0$ or $(|Q_L| + |W|)/T_H \geq |Q_L|/T_L$
 - Divide by $|W|$ and rearrange to get
 - $\beta = |Q_L| / |W| \leq T_L / (T_H - T_L) = \beta_{\text{Carnot}}$
- It is impossible to transfer heat from low to high temperature without work input

This Week's Homework

- Generally deal with definitions of efficiency, $\eta = |W| / |Q_H|$, or coefficient of performance, $\beta = |Q_L| / |W|$
- Also require use of $|Q_H| = |Q_L| + |W|$
- For Carnot Cycles, $\eta = 1 - T_L / T_H$
- A heat pump cools outside air while heating inside air; $\beta_{\text{heat pump}} = |Q_H| / |W|$

Class Exercise

- Find the entropy and enthalpy of water at a pressure of 1000 psia and 1000°F
- What is the enthalpy at a state with the same entropy you just found and a pressure of 80 psia?
- What is the enthalpy at a state with the same entropy you found above and a pressure of 5 psia?

Class Exercise Solution

- At 1000 psia and 1000°F, $h = 1506.2$ Btu/lb_m, $s = 1.6535$ Btu/lb_m·R (Table A-6E, p946)
- At 80 psia, $s = 1.6535$ Btu/lb_m·R is between $s = 1.6271$ Btu/lb_m·R at 320°F and $s = 1.6545$ Btu/lb_m·R at 360°F

Interpolate h and s

$$h = \frac{1187.9 \text{ kJ}}{\text{kg}} + \frac{\frac{1209.9 \text{ kJ}}{\text{kg}} - \frac{1187.9 \text{ kJ}}{\text{kg}}}{\frac{1.6545 \text{ kJ}}{\text{kg}} - \frac{1.6271 \text{ kJ}}{\text{kg}}} \cdot \left(\frac{1.6535 \text{ kJ}}{\text{kg}} - \frac{1.6271 \text{ kJ}}{\text{kg}} \right)$$

Class Exercise Solution II

- At 5 psia, $s = 1.6535$ Btu/lb_m·R is between $s_f = 0.23488$ Btu/lb_m·R and $s_g = 1.8438$ Btu/lb_m·R (Table A-5E, p 942)
- Compute $h = h_f + xh_{fg}$ and $x = (s - s_f)/s_{fg}$
 - Combine these equations to eliminate x

$$h = h_f + \frac{s - s_f}{s_{fg}} h_{fg} = \frac{130.18 \text{ kJ}}{\text{kg}} + \frac{\frac{1.6535 \text{ kJ}}{\text{kg}} - \frac{0.23488 \text{ kJ}}{\text{kg}}}{\frac{1.60894 \text{ kJ}}{\text{kg}}} \frac{1000.5 \text{ kJ}}{\text{kg}} = \frac{1012.3 \text{ kJ}}{\text{kg}}$$

Meaning of Class Exercise

- Entropy is a property
- We can find the entropy if we know two independent properties that can specify the state
- We can use entropy and one other property to identify the state
- Mixed region calculations for entropy are the same as those for volume, internal energy and enthalpy